

Discussion of “Direct Solutions for Uniform Flow Parameters of Wide Rectangular and Triangular Sections” by Ahmed A. Lamri, Said M. Easa, Mohamed T. Bouziane, Mohammad Bijankhan, and Yan-Cheng Han

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[https://doi.org/10.1061/\(ASCE\)IR.1943-4774.0001583](https://doi.org/10.1061/(ASCE)IR.1943-4774.0001583)

The authors have proposed new explicit equations for the head loss and normal depth of wide rectangular and triangular open channel sections using the Colebrook white equation and Lagrange’s inversion theorem.

The ASCE Task Force on friction factors in open channels (ASCE 1963) expressed its belief in the utility of using the Darcy-Weisbach formulation for resistance to flow in open channels for other than fully rough flow; however, these recommendations have almost been ignored, and the Manning-Strickler formulation continues to dominate, maybe because most pipes and channels flow under turbulent conditions (Boukhari et al. 2021; Zeghadnia 2007, 2014; Zeghadnia et al. 2014; Zeghadnia and Robert 2017).

Different mistakes were picked up through a deep analysis of the results; for instance, in the section “Modeling Friction Factor” of the original paper, the term for the hydraulic radius, R_h , was missed in Eqs. (8) and (10). They should be written as follows:

$$f = \left[c - 2.03 \log \left(\frac{\varepsilon R}{3.7088 R_h} + \frac{1}{\sqrt{f}} \right) \right]^{-2} \quad (1)$$

$$f^{1/2} = c + \sum_{n=1}^{\infty} \frac{\left(\frac{-2.03}{\ln 10}\right)^n}{\Gamma(n+1)} \frac{d^{n-1}}{dc^{n-1}} \left\{ \left[\frac{\varepsilon R}{3.7088 R_h} + c \right]^n \right\} \quad (2)$$

Eq. (11) of the original paper can be simplified and rewritten as follows:

$$f^{-1/2} = c + 2.03 \log \left[-1 + \frac{2.03}{b \ln 10} + 0.5 \left(\frac{2.03}{b \ln 10} \right)^2 \ln \left(\frac{b}{e^2} \right) + \left(\frac{2.03}{b \ln 10} \right)^3 \left(\frac{(\ln b)^2}{3} - \frac{3}{2} b + 1 \right) \right] \quad (3)$$

To estimate the accuracy of the proposed model, the authors have not taken enough values of ε/R_h and R to ensure the study of all possibilities; at first glance, it is easy to conclude that the computation steps were not small enough to judge the reliability of the proposed formulas. A deep investigation of the maximum error leads to the following:

To establish Fig. 1, random values from the ranges $10^{-6} \leq \varepsilon/R_h \leq 10^2$ and $4,000 \leq R \leq 10^8$ were used.

Fig. 1 shows the maximum error that can be induced when using Eqs. (11) and (14) of the original paper, where the maximum error of Eq. (11) is 0.16%; however, for Eq. (14), the maximum error is 0.141%, which is different from what the authors indicated.

In the section “Modeling Normal Depths” (rectangular form) of the original paper, the term y_n^* was missed in Eq. (21). So, Eq. (21) is given as

$$h = \frac{8\sqrt{2}\varepsilon^*}{37.088v^*y^*} \quad (4)$$

More than one million—random—values were used to check the reliability of Eq. (22); the results plotted in Fig. 2 shows, without doubt, a huge deviation between Eq. (18) of the original paper and the proposed formula (Fig. 2), where the maximum relative errors recorded was 6.723%.

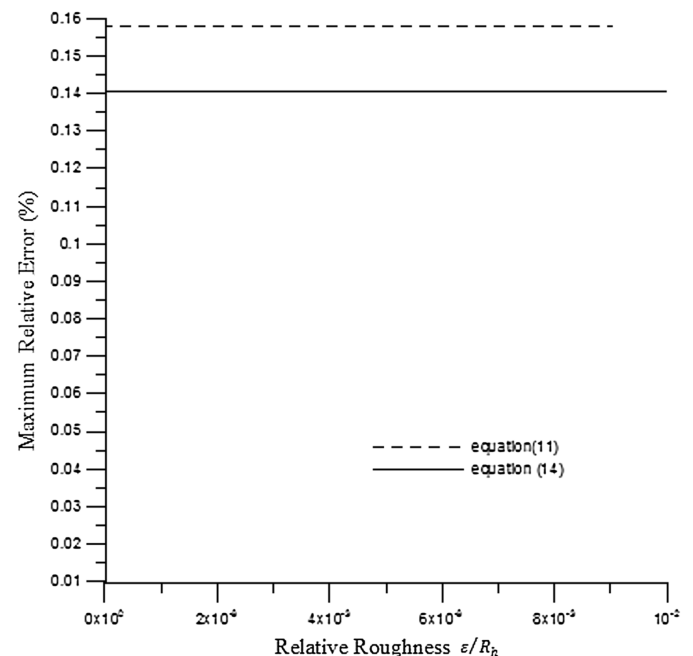


Fig. 1. The maximum deviation (in percent) for the entire range of: $10^{-6} \leq \varepsilon/R_h \leq 10^2$ and $4,000 \leq R \leq 10^8$ of Eqs. (11) and (14).

The maximum error recorded for Eq. (6) was equal to or less than 0.1% as shown Fig. 5.

It can be seen that the proposed equations do not reflect the accuracy provided by the authors, especially Eq. (22) of the original paper. Corrections and improvements are proposed to enhance the results such as Eqs. (5) and (6).

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